

3. Quantum Mechanics (Fall 2006)

Consider two flavours of massive neutrinos, denote $|\nu_e\rangle$ the electron neutrino flavour eigenstate and $|\nu_\mu\rangle$ the muon neutrino flavour eigenstate. These are related to the energy eigenstates $|\nu_1\rangle$ and $|\nu_2\rangle$ by

$$\begin{aligned} |\nu_e\rangle &= \cos(\theta) |\nu_1\rangle - \sin(\theta) |\nu_2\rangle \\ |\nu_\mu\rangle &= \sin(\theta) |\nu_1\rangle + \cos(\theta) |\nu_2\rangle \end{aligned}$$

- (a) Show that flavour eigenstates and energy eigenstates are related by a unitary transformation.
 (b) The energy of the eigenstate $|\nu_i\rangle$ is

$$E_i = \sqrt{\mathbf{p}^2 c^2 + m_i^2 c^4}, \quad i = 1, 2$$

Assume that an electron neutrino is produced in the sun with momentum \mathbf{p} such that $|\mathbf{p}| \gg m_i c$. Find the probability for the electron neutrino to oscillate into a muon neutrino after travelling a distance L .

$$\begin{aligned} \text{a) } \vec{\Psi}_f &= \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \vec{\Psi}_e & |\Psi\rangle &= \vec{\Psi}_f \cdot |\vec{\beta}_f\rangle = \psi_{f1} |\nu_e\rangle + \psi_{f2} |\nu_\mu\rangle \\ &= U \vec{\Psi}_e & &= \vec{\Psi}_e \cdot |\vec{\beta}_e\rangle = \psi_{e1} |\nu_1\rangle + \psi_{e2} |\nu_2\rangle \end{aligned}$$

$$\text{where, e.g., } \vec{\Psi}_f = \begin{pmatrix} \psi_{f1} \\ \psi_{f2} \end{pmatrix} \quad |\vec{\beta}_f\rangle = \begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix}$$

$$U^\dagger = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \Rightarrow U^\dagger U = \begin{pmatrix} c^2 + s^2 & -sc + cs \\ -cs + sc & -(s^2) + c^2 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = 1$$

$$\Rightarrow U^\dagger = U^{-1} \Rightarrow U \text{ is unitary}$$

$$\begin{aligned} \text{b) } |\Psi_0\rangle &= |\nu_e\rangle \text{ with } |\mathbf{p}| \gg m_1 c, m_2 c, \text{ traveling } L \Rightarrow \Delta t \approx \frac{L}{c} \equiv \tau \\ |\Psi(t)\rangle &= e^{-iHt/\hbar} |\Psi_0\rangle = e^{-i\omega_1 t} \cos \theta |\nu_1\rangle - e^{-i\omega_2 t} \sin \theta |\nu_2\rangle \text{ where } \omega_i \equiv \frac{E_i}{\hbar} \end{aligned}$$

$$\begin{aligned} P_{\nu_e \rightarrow \nu_\mu}(L) &= |\langle \nu_\mu | \Psi(\Delta t) \rangle|^2 \\ \langle \nu_\mu | \Psi(t) \rangle &= (\sin \theta \cos \theta) \begin{pmatrix} e^{-i\omega_1 t} \cos \theta \\ -e^{-i\omega_2 t} \sin \theta \end{pmatrix} = (e^{-i\omega_1 t} - e^{-i\omega_2 t}) \sin \theta \cos \theta \\ &= \frac{1}{2} \sin 2\theta (e^{-i\omega_1 t} - e^{-i\omega_2 t}) \end{aligned}$$

$$\begin{aligned} P_{\nu_e \rightarrow \nu_\mu}(L) &= \frac{1}{4} \sin^2 2\theta [1 - e^{i(\omega_2 - \omega_1)\tau} - e^{-i(\omega_2 - \omega_1)\tau} + 1] = \frac{1}{4} \sin^2 2\theta [2 - 2 \cos(\omega_2 - \omega_1)\tau] \\ &= \sin^2 2\theta \left(\frac{1}{2} - \frac{1}{2} \cos(\omega_2 - \omega_1)\tau \right) = \sin^2 2\theta \sin^2 \left(\frac{1}{2}(\omega_2 - \omega_1)\tau \right) \end{aligned}$$

$$\begin{aligned} \left(E_2 - E_1 = \sqrt{p^2 c^2 + m_2^2 c^4} - \sqrt{p^2 c^2 + m_1^2 c^4} = pc \left[(1 + m_2^2 c^2/p^2)^{1/2} - (1 + m_1^2 c^2/p^2)^{1/2} \right] \frac{m_i c}{p} \ll 1 \right) \\ \approx pc \left[1 + \frac{1}{2} m_2^2 c^2/p^2 - 1 - \frac{1}{2} m_1^2 c^2/p^2 \right] = c^3 \Delta(m^2)/2p \\ \approx \sin^2 2\theta \sin^2 [c^2 \Delta(m^2)L/4\hbar p] \end{aligned}$$