

### 3. Quantum Mechanics (Fall 2004)

A positron has the same mass  $m$  as the electron, but the opposite charge. Consider a set of states containing one electron and one positron. A complete set of these states can be labeled  $|\mathbf{r}_+, \mathbf{r}_-\rangle$ , where  $\mathbf{r}_+$  and  $\mathbf{r}_-$  are the positions of the positron and electron, respectively. Normalize these states so that

$$\langle \mathbf{r}_+, \mathbf{r}_- | \mathbf{r}'_+, \mathbf{r}'_- \rangle = \delta_3(\mathbf{r}'_+ - \mathbf{r}_+) \delta_3(\mathbf{r}'_- - \mathbf{r}_-)$$

Then if the system is in any state  $|\psi\rangle$ , the wave function is

$$\psi(\mathbf{r}_+, \mathbf{r}_-) = \langle \mathbf{r}_+, \mathbf{r}_- | \psi \rangle$$

In this problem ignore spin.

- In terms of  $\psi(\mathbf{r}_+, \mathbf{r}_-)$ , what is the probability that at least one of the two particles is farther than a distance  $b$  from the origin?
- Write down the Hamiltonian for this electron-positron system, including the electrostatic (Coulomb) interactions between the two particles.
- Let  $\mathbf{r} = \mathbf{r}_+ - \mathbf{r}_-$  and  $\mathbf{R} = \frac{1}{2}(\mathbf{r}_+ + \mathbf{r}_-)$ . Write the Hamiltonian in terms of the new coordinates and their canonically conjugate momenta  $\mathbf{p}$  and  $\mathbf{P}$ .
- The bound electron-positron system is called *positronium*. For states with zero total momentum, write a formula for the possible negative values of the energy<sup>1</sup>. What is the approximate numerical value, in electron volts, of the ground state energy?
- Define the *charge conjugation* operator  $C$  on this system by

$$C|\mathbf{r}_+, \mathbf{r}_-\rangle = |\mathbf{r}_-, \mathbf{r}_+\rangle$$

Show that  $C$  commutes with the Hamiltonian. What is the eigenvalue of  $C$  on the state of lowest energy?

a) At least one particle farther than  $r=b$

$\Rightarrow$  Not (both particles closer than  $r=b$ )  $V_b \equiv$  sphere, radius  $b$ , at origin

$$P = 1 - \int_{V_b} \int_{V_b} |\Psi(\vec{r}_+, \vec{r}_-)|^2 d\Omega_+ d\Omega_- = 1 - \int \int_{\text{all angles}} \int_0^b \int_0^b |\Psi(\vec{r}_+, \vec{r}_-)|^2 r_+^2 dr_+ r_-^2 dr_- d\Omega_+ d\Omega_-$$

b)  $H = \frac{\vec{p}_+^2}{2m} + \frac{\vec{p}_-^2}{2m} - K \frac{e^2}{|\vec{r}_+ - \vec{r}_-|}$

c)  $H = \frac{\vec{P}^2}{2M} + \frac{\vec{p}^2}{2\mu} - K \frac{e^2}{r} \quad M = 2m, \mu = \frac{1}{2}m$

d)  $E_n = -\frac{\mu K^2 e^4}{2\hbar^2 n^2} = -\frac{m K^2 e^4}{4\hbar^2} \frac{1}{n^2} \approx -\frac{1}{2} (13.6 \text{ eV}) \frac{1}{n^2} \quad E_1 \approx (6.8 \text{ eV})$

e) Since  $\hat{C}^2 = \hat{1}$ ,  $\hat{C}\hat{H}\hat{C} = \hat{H} \Leftrightarrow [\hat{C}, \hat{H}] = \hat{0}$

$$\begin{aligned} \hat{C}\hat{r}\hat{C}|\vec{r}_+, \vec{r}_-\rangle &= \hat{C}\hat{r}|\vec{r}_-, \vec{r}_+\rangle = \hat{C}(\hat{r}_+ - \hat{r}_-)|\vec{r}_-, \vec{r}_+\rangle = \hat{C}(\vec{r}_- - \vec{r}_+)|\vec{r}_-, \vec{r}_+\rangle \\ &= (\vec{r}_- - \vec{r}_+)\hat{C}|\vec{r}_-, \vec{r}_+\rangle = -(\vec{r}_+ - \vec{r}_-)|\vec{r}_+, \vec{r}_-\rangle = -\hat{r}|\vec{r}_+, \vec{r}_-\rangle \end{aligned}$$

$\hat{C}\hat{r}\hat{C} = -\hat{r}$  for a complete basis  $\Rightarrow \hat{C}\hat{r}\hat{C} = -\hat{r}$  (for all state kets)

<sup>1</sup>Write your answer in terms of  $m$ ,  $e^2$  or  $\alpha$ ,  $\hbar$ ,  $c$ , the Bohr radius, etc. You may use units in which  $\hbar = c = 1$ .

### 3. Quantum Mechanics (Fall 2004)

e) (continued)

$$\begin{aligned}\hat{C}\hat{p}\hat{C}|\vec{r}_+, \vec{r}_-\rangle &= \hat{C}(\hat{p}_+ - \hat{p}_-)|\vec{r}_-, \vec{r}_+\rangle = \hat{C}(-i\hbar)(\vec{\nabla}_- - \vec{\nabla}_+)|\vec{r}_-, \vec{r}_+\rangle \\ &= (-i\hbar)(\vec{\nabla}_- - \vec{\nabla}_+)\hat{C}|\vec{r}_-, \vec{r}_+\rangle = -(-i\hbar)(\vec{\nabla}_+ - \vec{\nabla}_-)|\vec{r}_+, \vec{r}_-\rangle \\ &= -(\hat{p}_+ - \hat{p}_-)|\vec{r}_+, \vec{r}_-\rangle = -\hat{p}|\vec{r}_+, \vec{r}_-\rangle \quad \Rightarrow \hat{C}\hat{p}\hat{C} = -\hat{p}\end{aligned}$$

$$\hat{C}\hat{p}^2\hat{C} = \hat{C}\frac{1}{2}(\hat{p}_+ + \hat{p}_-)\hat{C} = \frac{1}{2}(\hat{p}_- + \hat{p}_+) = \hat{p}^2 \quad \Rightarrow \hat{C}\hat{p}^2\hat{C} = \hat{p}^2$$

$$\begin{aligned}\Rightarrow \hat{C}\hat{H}\hat{C} &= \hat{C}\left[\frac{\hat{p}^2}{2M} + \frac{\hat{p}^2}{2\mu} - \frac{Ke^2}{|\vec{r}|}\right]\hat{C} = \left[\frac{(\hat{p})^2}{2M} + \frac{(-\hat{p})^2}{2\mu} - \frac{Ke^2}{|\vec{r}|}\right] \\ &= \left[\frac{\hat{p}^2}{2M} + \frac{\hat{p}^2}{2\mu} - \frac{Ke^2}{|\vec{r}|}\right] = \hat{H}\end{aligned}$$

$$\Rightarrow [\hat{C}, \hat{H}] = \hat{0} \quad \checkmark$$

Since the lowest energy eigenstate is spherically symmetric with respect to  $\vec{r}$  and  $\hat{C}$  acts as the spatial parity operator with respect to  $\vec{r}$ , the lowest energy eigenstate is even in  $\vec{r}$  and is an eigenstate of  $\hat{C}$  with eigenvalue  $+1$ .